**Electrostatic Problems with Dielectrics**

$$
C = \frac{Q}{V} = \frac{A}{d_2 \left(\frac{A}{\varepsilon_2} + \frac{d_1}{\varepsilon_1}\right)}
$$

## **Examples of Electrostatic Problems with Dielectrics**

Problem: Find  $\vec{D}$  (electric flux density),  $\vec{E}$  (electric field intensity), and (polarization) for a metallic sphere (radius *a*, charge *Q*), coated by a dielectric (radius *b*), and the charge densities at the interfaces.  $\vec{D}$  (electric flux density),  $\vec{E}$  (electric field intensity), and  $\vec{P}$ 

Solution: Use Gauss' Law  $\oint \vec{D} \cdot d\vec{s} = \int \rho_{\psi} dv$ *s v*  $\oint_{s} \vec{D} \cdot d\vec{s} = \int_{v} \rho$ 

In  $region \ 0, \ r < a, \vec{D} = 0, \vec{E} = 0$ 

In *region* 1,  $a < r < b$ :

$$
\vec{D}_1 = \hat{e}_r \frac{Q}{4\pi r^2}
$$

In *region* 2,  $r > b$ :

$$
\vec{D}_2 = \hat{e}_r \frac{Q}{4\pi r^2}
$$



The electric fields are:

In *region* 1,  $a < r < b$ :

$$
\vec{E}_1 = \hat{e}_r \frac{Q}{4\pi \varepsilon_1 r^2}
$$

In *region* 2, *r > b*:

$$
\vec{E}_2 = \hat{e}_r \frac{Q}{4\pi \varepsilon_o r^2}
$$

Note that the electric field varies with the medium, and the flux density does not.

The polarization is:

$$
\vec{P} = \varepsilon_{o} \chi_{e} \vec{E} \text{ and } \varepsilon_{r} = 1 + \chi_{e}
$$

In *region* 1:

$$
\vec{P}_1 = \varepsilon_o \chi_e \vec{E}_1 = \varepsilon_o (\varepsilon_r - 1) \vec{E}_1 = \hat{e}_r \frac{\varepsilon_o (\varepsilon_r - 1) Q}{4 \pi \varepsilon_1 r^2}
$$

#### In *region* 2:  $\rightarrow$  $P_2 = 0$

Now the surface charge on  $0-1$  boundary. The boundary condition is:

$$
\vec{n} \cdot (\vec{P}_1 - \vec{P}_o) = -\rho_{ps},
$$
  
or  $\hat{e}_r \cdot \left[ \hat{e}_r \frac{\varepsilon_o (\varepsilon_r - 1) Q}{4\pi \varepsilon_1 a^2} - 1 \right] = -\rho_{ps}$ 

Then:

$$
\left. \rho_{ps} \right|_{r=a} = -\frac{\varepsilon_{o} (\varepsilon_{r} - 1) Q}{4 \pi \varepsilon_{1} a^{2}}
$$

Now for  $r = b$ :

$$
\hat{e}_r \cdot \left[0 - \hat{e}_r \frac{\varepsilon_o(\varepsilon_r - 1)Q}{4\pi \varepsilon_1 b^2}\right] = -\rho_{ps}
$$

Then:



(a) The electric flux density in all regions where it may be seen that  $\overrightarrow{D}$  is continuous. region 2 and hence was presented by closer flux lines. (c) The polarization  $\vec{P}$  that (b) The electric field intensity  $\vec{E}$  in all regions where it may be seen that  $\vec{E}$  is larger in (b) The electric field intensity  $\vec{E}$  in all regions where it may be seen that  $\vec{E}$  is larger in only exists in region 1 where we have dielectric material and (d) the polarization surface charge on both surfaces of region 1.

## **Example**

Parallel plate capacitor with two different dielectrics:



Determine (1) the electric field intensities and the electric flux density vectors, and (2) the surface charge density of polarization charges.

(1)  $\varepsilon_1 = \varepsilon_o \varepsilon_{r1}$ 

The electric field is:  $V = \int_{a}^{d} \vec{E} \cdot d\vec{l} = E d$  or  $E = V/d$  , everywhere  $\vec{D}_1 = \varepsilon_1 \vec{E}_1$  (*region* 1) and  $\vec{D}_2 = \varepsilon_2 \vec{E}_2$  (*region* 2)  $=\int_0^d \vec{E} \cdot d\vec{l}$  = 0

Or, in terms of *V*:

$$
\vec{D}_1 = \varepsilon_1 \frac{V}{d}
$$
 and  $\vec{D}_2 = \varepsilon_2 \frac{V}{d}$ 

(2) At the surface of the conductor, the surface free charge density is:

$$
\rho_s = D_n
$$

Therefore, in *region* 1: and in *region* 2:  $\rho_{s2} = D_{n2} = \varepsilon_2 \frac{V}{d}$ *- - - - - - - - - - - -*  $\rho_{s1} = D_{n1} = \varepsilon_1 \frac{V}{d}$   $E \begin{bmatrix} \varepsilon_1 \\ \varepsilon_2 \end{bmatrix}$   $\varepsilon_1$  $\rho_{s1}$   $\varepsilon_1 < \varepsilon_2$   $\rho_{s2}$  $\rho_{s1} < \rho_{s2}$ 

The surface charge density of polarization charge is:

$$
-\rho_{ps1} = p_{normal1} \rightarrow \text{since } \vec{P} = \varepsilon_o \chi_e \vec{E} \text{ and } \varepsilon_{r1} = 1 + \chi_{e1}
$$
  

$$
\Rightarrow \vec{P}_1 = \varepsilon_o (\varepsilon_{r1} - 1) \vec{E}_1, \text{ or } \vec{P}_1 = (\varepsilon_1 - \varepsilon_o) \frac{V}{d} = -\rho_{ps1}
$$

and for *region* 2:

$$
\Rightarrow \vec{P}_2 = \varepsilon_o \left( \varepsilon_{r2} - 1 \right) \vec{E}_1 \text{, or } \vec{P}_2 = \left( \varepsilon_{2} - \varepsilon_o \right) \frac{V}{d} = -\rho_{ps2}
$$



Free surface charge density distribution  $\rho_s$  and the polarization charge density distribution  $\rho_{ps}$  on the interface between the conducting planes and the dielectric materials. The total charge density on the surface of the conductor is constant and equal to  $\varepsilon_o E$ .

The total charge density (sum of bound and free) is a constant  $=\boldsymbol{\mathcal{E}}_{o}\,E$  .

## **Capacitance** (We need this for transmission lines!)

A measure of the amount of charge a particular arrangement of two conducfors is able to retain per unit voltage:  $C \equiv Q/V$  .

Example: Spherical capacitor

$$
\oint_{s} \vec{D} \cdot d\vec{s} = \int_{v} \rho_{v} dv
$$

Since  $\vec{D}$  =  $\varepsilon \,\vec{E}$  and the total charge on the inner sphere is  $Q$ , we have:  $\rightarrow$  $D = \varepsilon E$ 

$$
\vec{E} = \frac{Q\hat{e}_r}{4\pi\,\varepsilon\,r^2}
$$

and the potential difference is:  $V_{ab}\!=\!-\!\int_b^a\vec{E}\!\cdot d\vec{l}=\int_a^b\vec{E}\!\cdot d\vec{l}$ *b a*  $=-\int_b^a \vec{E} \cdot d\vec{l} = \int_a^b \vec{E} \cdot d\vec{l}$ 

Discussion: The minus sign indicates that positive work must be done to move a positive test charge from  $b$  to  $a$  against the field direction (radially out of  $a$  to  $b$ ). Now, because of symmetry,  $d\vec{l} = dr \hat{e}_r$ .

Then,

$$
V_{ab} = -\int_{a}^{b} \left[ \frac{Q(\hat{e}_{r} \cdot \hat{e}_{r})}{4 \pi \epsilon r^{2}} \right] dr = \left[ -\frac{Q}{4 \pi \epsilon r} \right]_{a}^{b}
$$

$$
= \left( -\frac{Q}{4 \pi \epsilon} \right) \left( -\frac{1}{a} + \frac{1}{b} \right) = \frac{Q}{4 \pi \epsilon} \left( \frac{1}{a} - \frac{1}{b} \right)
$$

The capacitance is:

$$
C \equiv \frac{Q}{V} = \frac{4\pi\,\varepsilon}{\frac{1}{b} - \frac{1}{a}}
$$

Another example: Infinitely wide parallel plate capacitor. Find capacitance/unit area.



There are two fields,  $E^{\phantom{\dagger}}_1$  and  $E^{\phantom{\dagger}}_2$ , both directed normal to the plates.

Gauss' Law gives: 
$$
E_1 = \frac{\rho_s}{\varepsilon_1}
$$
 and  $E_2 = \frac{\rho_s}{\varepsilon_2}$ .

The flux density  $\dot{D}$  is (and must be) continuous across the dielectric boundary, and its value equals the surface charge density.  $\rightarrow$ *D*

The potential is given by:

$$
V = -\int_{d_2}^{0} E_2 \cdot d\vec{l} - \int_{d_1 + d_2}^{d_2} E_1 \cdot d\vec{l} = \frac{\rho_s d_2}{\varepsilon_2} + \frac{\rho_s d_1}{\varepsilon_1}
$$
  
=  $\rho_s \left( \frac{d_2}{\varepsilon_2} + \frac{d_1}{\varepsilon_1} \right)$ 

Now  $Q = \rho_s A$ .

So the capacitance is:  $C = \frac{Q}{V}$ *V A*  $=\frac{Q}{V}=\frac{A}{d_2 \bigg / \bigg / \bigg / \bigg / \bigg / \bigg }$  $^{2}/_{c}$  + 2 1  $\left[{\varepsilon} _2\right.^\top\left/ {\varepsilon} _1\right]$ 

# **Another important example:**



Again, Gauss' Law:

$$
\oint_{s} \mathcal{E}_{1} \vec{E} \cdot d\vec{s} = \int_{v} \rho_{v} dv
$$

or for  $a < r < b$ :  $\varepsilon_1 E 2\pi r L = \rho_s 2\pi a L$ , or  $E_r = \frac{\rho_s a}{\varepsilon_1 r}$ 

The voltage is 
$$
V = -\int_b^a E \cdot d\vec{l} = -\frac{\rho_s a}{\varepsilon_1} \int_b^a \frac{dr}{r} = \frac{\rho_s a}{\varepsilon_1} \ln\left(\frac{b}{a}\right)
$$
,

so that  $C = \frac{2 \pi a \rho_s}{L} = \frac{\ln(b/a)}{L}$ . *V unit length s*  $=\frac{2\pi a \rho_s}{N} = \frac{\ln(b/a)}{1}$  $2\pi\varepsilon_1$  $\pi a \rho$  $\frac{\pi \, \varepsilon_1}{\ln \left(\frac{b}{a}\right)}$ 

**Another important example** Two conductor transmission line.



There are equal and opposite charges per unit length on the conductors. Let  $+r_l$  be on the right conductor. We will use the principal of superposition to solve this problem.

Gauss' Law for the right hand wire yields for the field:

$$
\int_{s} \varepsilon \vec{E}_{1} \cdot d\vec{s} = \int_{c} r_{l} dl
$$
\n
$$
\int_{\phi=0}^{2\pi} \int_{z=0}^{l} \varepsilon \vec{E}_{1r} \hat{e}_{r} \cdot r \, d\phi \, dz \, \hat{e}_{r} = r_{l} l
$$
\n
$$
E_{1r} = \frac{r_{l}}{2\pi \varepsilon r}
$$

The electric field at  $\,d_1$  due to right hand wire is

$$
\vec{E}_1 = \frac{r_l}{2\pi \varepsilon d_1} \hat{e}_r
$$

Now, the left hand wire:

$$
\vec{E}_2 = \frac{r_l(-\hat{e}_r)}{2\pi \varepsilon (d - d_1)}
$$

The total field is:

$$
\vec{E} = \vec{E}_1 + \vec{E}_2 = \left[ -\frac{r_l}{2\pi \varepsilon (d - d_1)} + \frac{r_l}{2\pi \varepsilon d_1} \right]
$$

$$
V = -\int_{d-a}^{a} \left[ \frac{r_l}{2\pi \varepsilon r} - \frac{r_l}{2\pi \varepsilon (d - r)} \right] \hat{e}_r \cdot dr \hat{e}_r = -\int_{d-a}^{a} \vec{E} \cdot d\vec{l}
$$

 $C=?$